

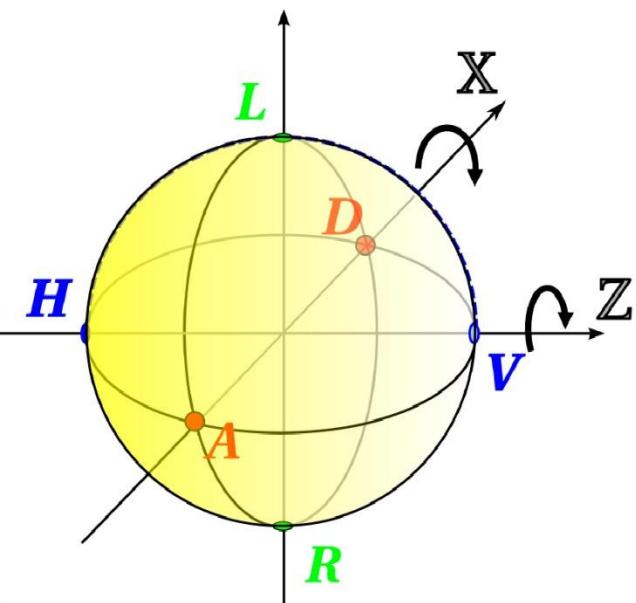
Jones vector by R. C. Jones in 1941

http://en.wikipedia.org/wiki/Jones_vector

$$\begin{pmatrix} E_x(t) \\ E_y(t) \end{pmatrix} = \begin{pmatrix} E_{0x} e^{i(kz - \omega t + \phi_x)} \\ E_{0y} e^{i(kz - \omega t + \phi_y)} \end{pmatrix} = \begin{pmatrix} E_{0x} e^{i\phi_x} \\ E_{0y} e^{i\phi_y} \end{pmatrix} e^{i(kz - \omega t)}$$

Polarization	Corresponding Jones vector	Typical ket Not
Linear polarized in the x-direction Typically called 'Horizontal'	$\begin{pmatrix} 1 \\ 0 \end{pmatrix}$	$ H\rangle$
Linear polarized in the y-direction Typically called 'Vertical'	$\begin{pmatrix} 0 \\ 1 \end{pmatrix}$	$ V\rangle$
Linear polarized at 45° from the x-axis Typically called 'Diagonal' L+45	$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$	$ D\rangle = \frac{1}{\sqrt{2}}(H\rangle + V\rangle)$
Linear polarized at -45° from the x-axis Typically called 'Anti-Diagonal' L-45	$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$	$ A\rangle = \frac{1}{\sqrt{2}}(H\rangle - V\rangle)$
Right Hand Circular Polarized Typically called RCP or RHCP	$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$	$ R\rangle = \frac{1}{\sqrt{2}}(H\rangle - i V\rangle)$
Left Hand Circular Polarized Typically called LCP or LHCP	$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ +i \end{pmatrix}$	$ L\rangle = \frac{1}{\sqrt{2}}(H\rangle + i V\rangle)$

Bloch sphere



Jones matrices

Optical element	Corresponding Jones matrix
Linear polarizer with axis of transmission horizontal ^[1]	$\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$
Linear polarizer with axis of transmission vertical ^[1]	$\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$
Linear polarizer with axis of transmission at $\pm 45^\circ$ with the horizontal ^[1]	$\frac{1}{2} \begin{pmatrix} 1 & \pm 1 \\ \pm 1 & 1 \end{pmatrix}$
Right circular polarizer ^[1]	$\frac{1}{2} \begin{pmatrix} 1 & i \\ -i & 1 \end{pmatrix}$
Left circular polarizer ^[1]	$\frac{1}{2} \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix}$
Linear polarizer with axis of transmission at angle θ with the horizontal. (Shown construction from rotating up from the horizontal into the polarizing element, the polarizing element, and then rotating back down into the horizontal.)	$\begin{pmatrix} \cos^2(\theta) & \cos(\theta)\sin(\theta) \\ \sin(\theta)\cos(\theta) & \sin^2(\theta) \end{pmatrix} =$ $\begin{pmatrix} \cos(\theta) & -\sin(\theta) \\ \sin(\theta) & \cos(\theta) \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} \cos(-\theta) & -\sin(-\theta) \\ \sin(-\theta) & \cos(-\theta) \end{pmatrix}$

Phase retarders	Corresponding Jones matrix
Quarter-wave plate with fast axis vertical ^{[2][note 1]}	$e^{i\pi/4} \begin{pmatrix} 1 & 0 \\ 0 & -i \end{pmatrix}$
Quarter-wave plate with fast axis horizontal ^[2]	$e^{i\pi/4} \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}$
Half-wave plate with fast axis at angle θ w.r.t the horizontal axis ^[3]	$\begin{pmatrix} \cos 2\theta & \sin 2\theta \\ \sin 2\theta & -\cos 2\theta \end{pmatrix}$
Any birefringent material (phase retarder) ^[4]	$\begin{pmatrix} e^{i\phi_x} \cos^2 \theta + e^{i\phi_y} \sin^2 \theta & (e^{i\phi_x} - e^{i\phi_y}) \cos \theta \sin \theta \\ (e^{i\phi_x} - e^{i\phi_y}) \cos \theta \sin \theta & e^{i\phi_x} \sin^2 \theta + e^{i\phi_y} \cos^2 \theta \end{pmatrix}$

Stokes parameters

by George Gabriel Stokes in 1852

http://en.wikipedia.org/wiki/Stokes_parameters

$$S_0 = I$$

$$S_1 = pI \cos 2\psi \cos 2\chi$$

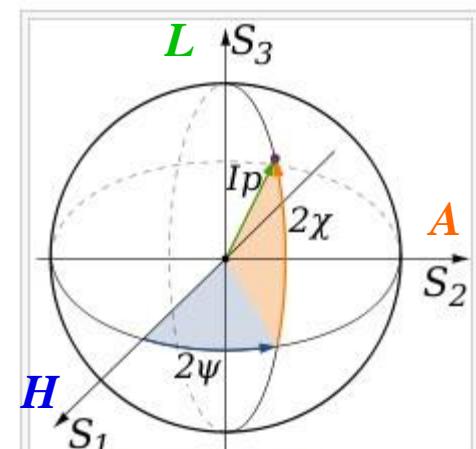
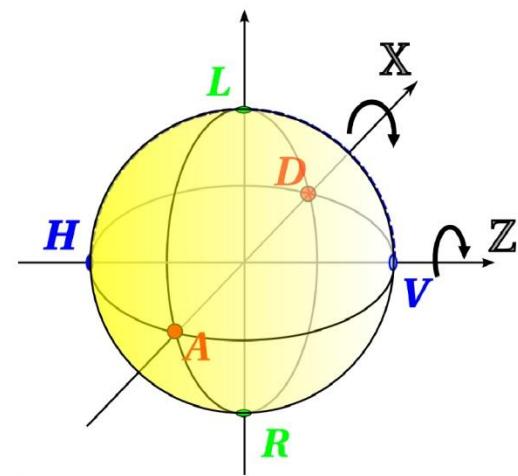
$$S_2 = pI \sin 2\psi \cos 2\chi$$

$$S_3 = pI \sin 2\chi$$

Stokes vectors

$$\vec{S} = \begin{pmatrix} S_0 \\ S_1 \\ S_2 \\ S_3 \end{pmatrix} = \begin{pmatrix} I \\ Q \\ U \\ V \end{pmatrix}$$

Bloch (Poincaré) sphere



The Poincaré sphere is the parametrisation of the last three Stokes' parameters in spherical coordinates

各种偏振态

$$S_{\text{LHP}} = I_0 \begin{pmatrix} 1 \\ 1 \\ 0 \\ 0 \end{pmatrix},$$



$$S_{\text{LVP}} = I_0 \begin{pmatrix} 1 \\ -1 \\ 0 \\ 0 \end{pmatrix},$$



$$S_{\text{L+45P}} = I_0 \begin{pmatrix} 1 \\ 0 \\ 1 \\ 0 \end{pmatrix},$$



$$S_{\text{L-45P}} = I_0 \begin{pmatrix} 1 \\ 0 \\ -1 \\ 0 \end{pmatrix}, \quad S_{\text{RCP}} = I_0 \begin{pmatrix} 1 \\ 0 \\ 0 \\ 1 \end{pmatrix},$$



$$S_{\text{LCP}} = I_0 \begin{pmatrix} 1 \\ 0 \\ 0 \\ -1 \end{pmatrix},$$



Stokes parameters

$$\begin{aligned} I &\equiv \langle E_x \rangle^2 + \langle E_y \rangle^2, \\ I &= \langle E_a \rangle^2 + \langle E_b \rangle^2, \\ I &= \langle E_l \rangle^2 + \langle E_r \rangle^2, \\ Q &\equiv \langle E_x \rangle^2 - \langle E_y \rangle^2, \\ U &\equiv \langle E_a \rangle^2 - \langle E_b \rangle^2, \\ V &\equiv \langle E_l \rangle^2 - \langle E_r \rangle^2, \end{aligned} \quad \vec{S} = \begin{pmatrix} S_0 \\ S_1 \\ S_2 \\ S_3 \end{pmatrix} = \begin{pmatrix} I \\ Q \\ U \\ V \end{pmatrix}$$

x, y 水平垂直线偏振 (**H,V**)

a, b +/-45° 偏振基 (**A,D**)

l, r 左右旋圆偏振基 (**L,R**)

Degree of polarization (DOP)

The DOP \mathcal{P} is defined by the equation

$$\mathcal{P} = \frac{I_{\text{pol}}}{I_{\text{tot}}} = \frac{\sqrt{S_1^2 + S_2^2 + S_3^2}}{S_0}, \quad 0 \leq \mathcal{P} \leq 1,$$

where I_{tot} is the total intensity. These results show that the relation between the Stokes parameters must be broadened to

$$S_0^2 \geq S_1^2 + S_2^2 + S_3^2,$$

where the $=$ and $>$ sign indicate completely and unpolarized/partially polarized light, respectively.